



# Resonant Oscillations of Liquid Marbles

Glen McHale\*, Stephen J. Elliott\*, Michael I. Newton\*,  
Dale L. Herbertson\* and Kadir Esmer\$

\*School of Science & Technology, Nottingham Trent University, UK

\$Department of Physics, Kocaeli University, Turkey

6<sup>th</sup> International Meeting on Electrowetting  
UCLA 20<sup>th</sup>-22<sup>nd</sup> August 2008

# Overview

## 1. Vibrations of Sessile Droplets

- Spherical and sessile droplets
- Type 1 and Type 2 modes
- Model vibrations

## 2. Liquid Marbles

- Concept: Perfect non-wetting droplets
- Mobility: Zero contact angle hysteresis
- Size dependence: Spherical to puddle

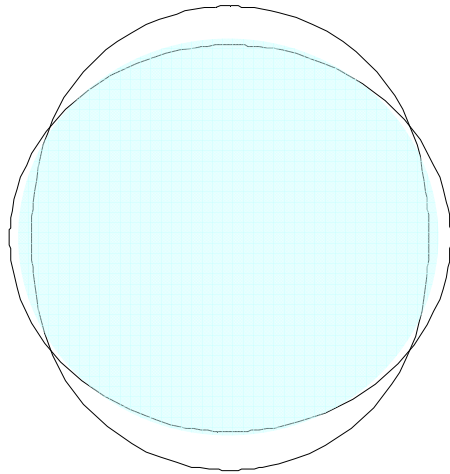
## 3. Oscillation Experiments

- Experimental set-up
- Spectra
- Volume and mode dependence

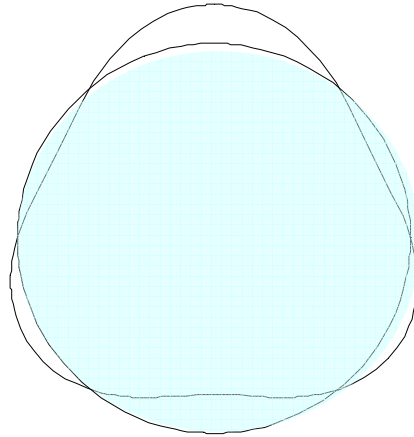
# Vibration of Sessile Droplets

# Vibrations of Spherical Droplets

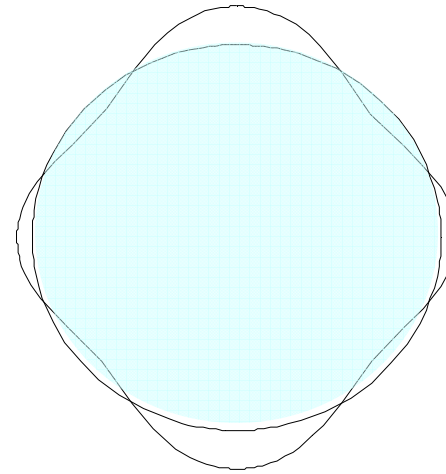
## Shape Modes



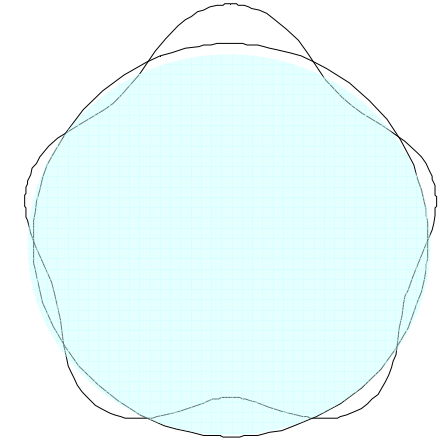
$n=2$



$n=3$



$n=4$



$n=5$

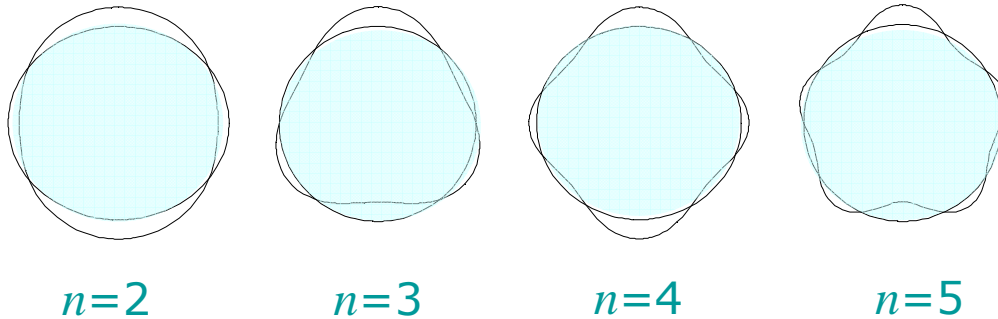
Number of nodes =  $2n$

## Resonant Frequencies

$$f_n^2 = \frac{n(n-1)(n+2)\gamma_{LV}}{4\rho\pi^2 R^3} \rightarrow \frac{n^3 \gamma_{LV}}{3\pi\rho V}$$

# Gravity-Capillary Waves

## Wavelengths



$n$  wavelengths fitted  
around perimeter  
(stationary surface waves)

$$p = 2\pi R = n\lambda_n$$

## 1-D Dispersion Equation

Bath of depth  $h$ , wave vector  $q = 2\pi/\lambda$

$\gamma_{LV}$  is surface tension and  $\rho$  is density

$\omega = 2\pi f$  is angular frequency

$$\omega^2 = \left( \underset{\substack{\uparrow \\ \text{gravity}}}{gq} + \frac{\underset{\substack{\uparrow \\ \text{surface tension}}}{\gamma_{LV}q^3}}{\underset{\substack{\uparrow \\ \text{finite depth}}}{\rho}} \right) \tanh(qh)$$

$$\omega_n^2 \approx \frac{\gamma_{LV} q_n^3}{\rho} = \frac{\gamma_{LV} (2\pi)^3}{\rho \lambda_n^3} \xrightarrow{2\pi R = n\lambda_n} f_n^2 \approx \frac{n^3 \gamma_{LV}}{4\rho\pi^2 R^3} = \frac{n^3 \gamma_{LV}}{3\pi\rho V}$$

# Sessile Droplets

## Effect of Substrate

1. Breaks complete spherical shape

Small droplets are spherical caps

Larger droplets flatten and have gravitational effects

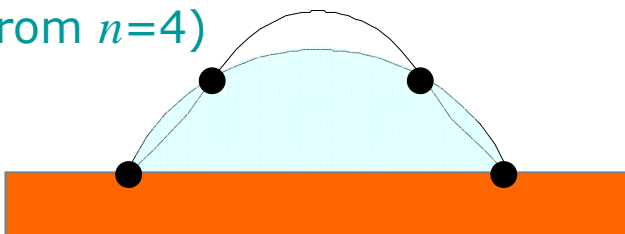
⇒ *Introduces dependence on contact angle,  $\theta$ , and volume,  $V$*

2. Contact line imposes boundary conditions

⇒ *Introduces dependence on contact angle hysteresis,  $\Delta\theta$*

### Type 1 – Immobile Contact Line

$j=2$  (from  $n=4$ )



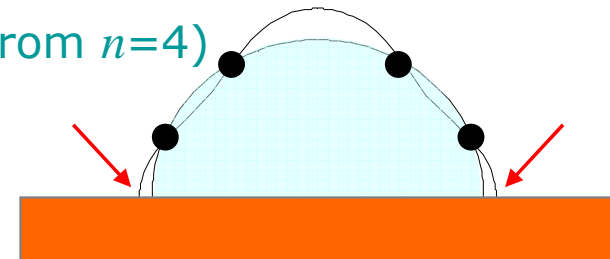
Number of nodes =  $2j$

Nodes at contact line

*Constant contact diameter*

### Type 2 – Mobile Contact Line

$k=2$  (from  $n=4$ )

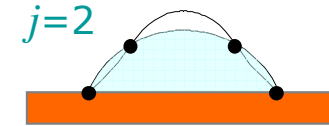


Number of nodes =  $2k$

Anti-nodes at contact line

*Constant contact angle,  $\theta_e$*

# Type 1 – Immobile Contact Line



## Pseudo-wavelength

Fit  $(j+1)$  half-wavelengths around perimeter length,  $p$ :  $2p = (j+1)\lambda_j$

Perimeter length depends on droplet volume and contact angle,  $p=p(V, \theta_e)$

Substitute into dispersion equation:

$$f_j^2 = \frac{(j+1)g}{4\pi p} \left( 1 + \left( \frac{(j+1)\pi}{p\kappa} \right)^2 \right) \tanh \left( \frac{(j+1)\pi h}{p} \right)$$

where  $\kappa^{-1} = (\gamma_{LV}/\rho g)^{1/2}$  is the capillary length

Noblin *et al* suggest mean height of droplet be used for depth,  $h_m = V/\pi R^2$

## Limits

Capillary regime:  $f \propto (j+1)^{3/2}$       3/2 power law for mode

Gravity regime:  $f \propto (j+1)$       linear dependence on mode

# Type 2 – Mobile Contact Line

## Pseudo-wavelength

Fit  $k$  wavelengths around perimeter length,  $p$ :

Perimeter length depends on droplet volume and contact angle,  $p=p(V, \theta_e)$

Substitute into dispersion equation:

$$f_k^2 = \frac{kg}{2\pi\rho} \left( 1 + \left( \frac{2\pi k}{p\kappa} \right)^2 \right) \tanh \left( \frac{2\pi h k}{p} \right)$$

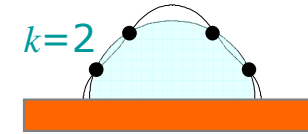
## Limits

Capillary regime:	$f \propto k^2$	3/2 power law for mode
Gravity regime:	$f \propto k$	linear dependence on mode

## Mixed modes

Pure type 2 is difficult to study

High amplitude excitation and contact angle hysteresis can lead to both contact line motion and contact angle oscillating simultaneously



$$p = k\lambda_k$$



# Liquid Marbles

# Experimental Design

## Requirements

1. For approximation to sphere, need very large ( $\gg 90^\circ$ ) contact angles
2. For type 2 mobile contact lines, need very low contact angle hysteresis

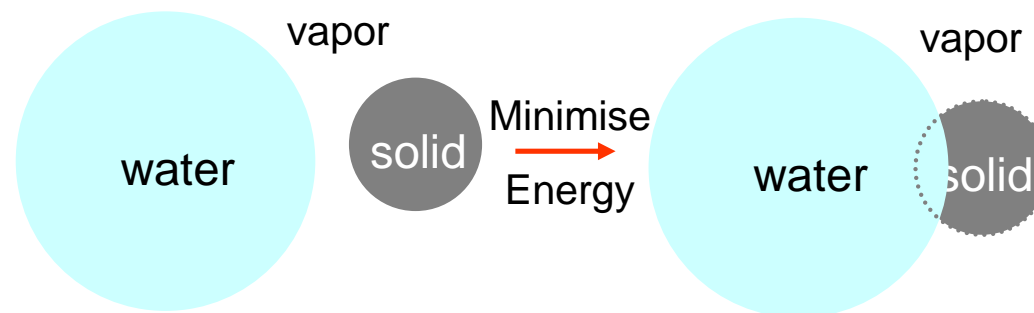
⇒ Lithographically fabricated superhydrophobic surface?

Problem in obtaining mobile contact lines

*Instigation of Cassie-Baxter to Wenzel transition*

OR use liquid marbles on Teflon AF1600/metal coated glass substrate

Surface free energy favors hydrophobic grains to attach to liquid-vapor interface



$$\Delta F = -\pi R_g^2 \gamma_{LV} (1 + \cos \theta_e)^2$$

Energy is always reduced on grain attachment

# Liquid Marbles as Superhydrophobic Surfaces

## Reversibility Idea

Make the solid “pillars” adhere more to the liquid than the substrate

Provides insulating “pillars” conformal to the liquid shape

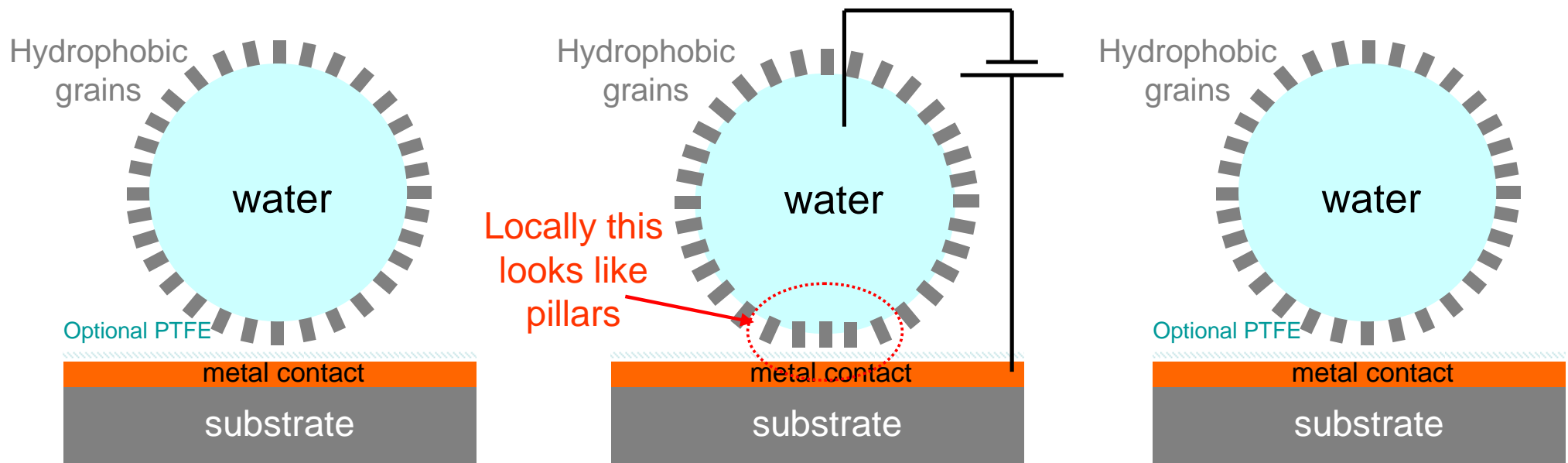
More hydrophobic grains “stick out” further (= to taller pillars)

Spin coated Teflon AF1600 on substrate to stop complete breakthrough if powder/granular coating is breached

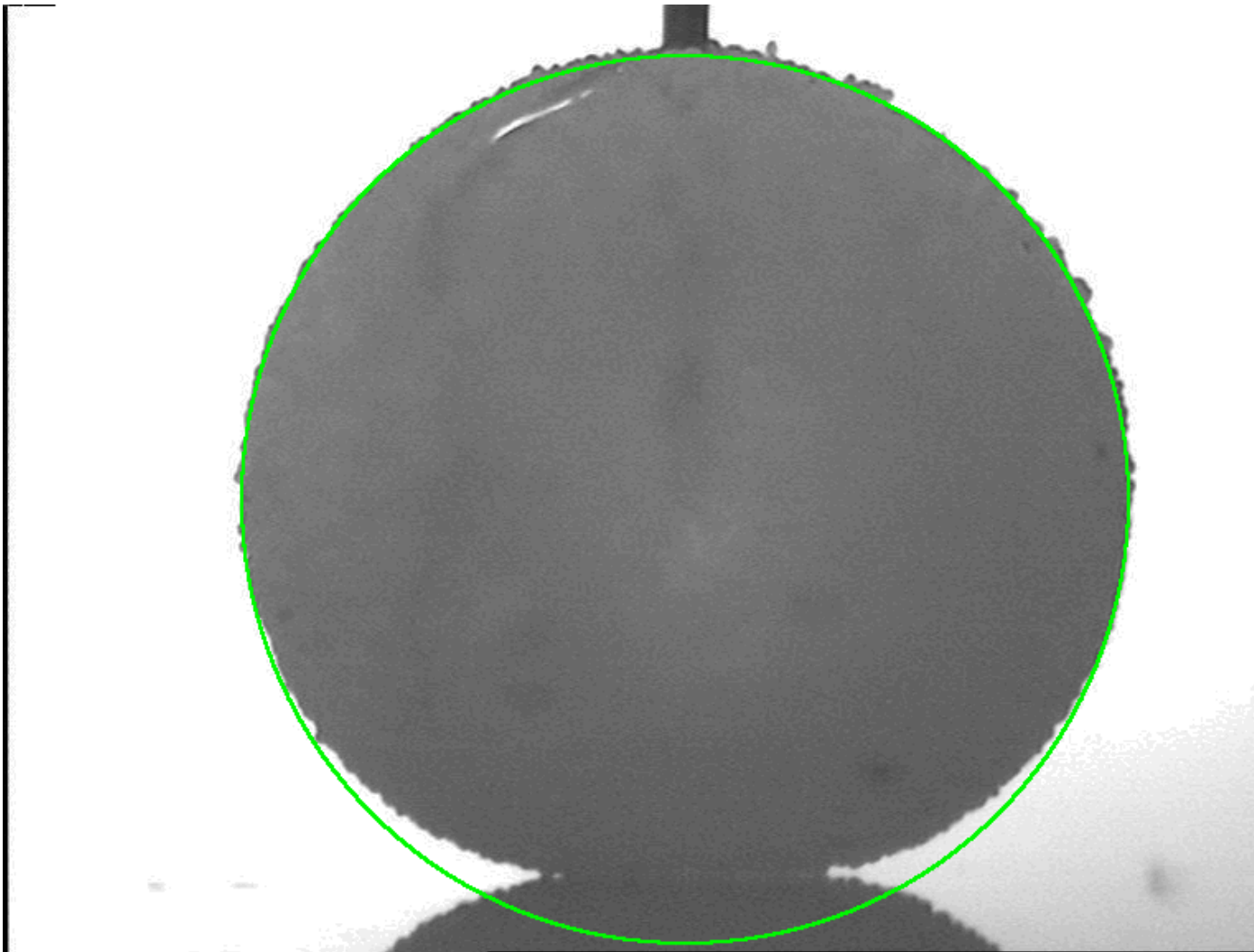
## Initial Shape

## Apply Voltage

## Remove Voltage



# Reversibility – Low V Cycle



# Mobility of Liquid Marbles



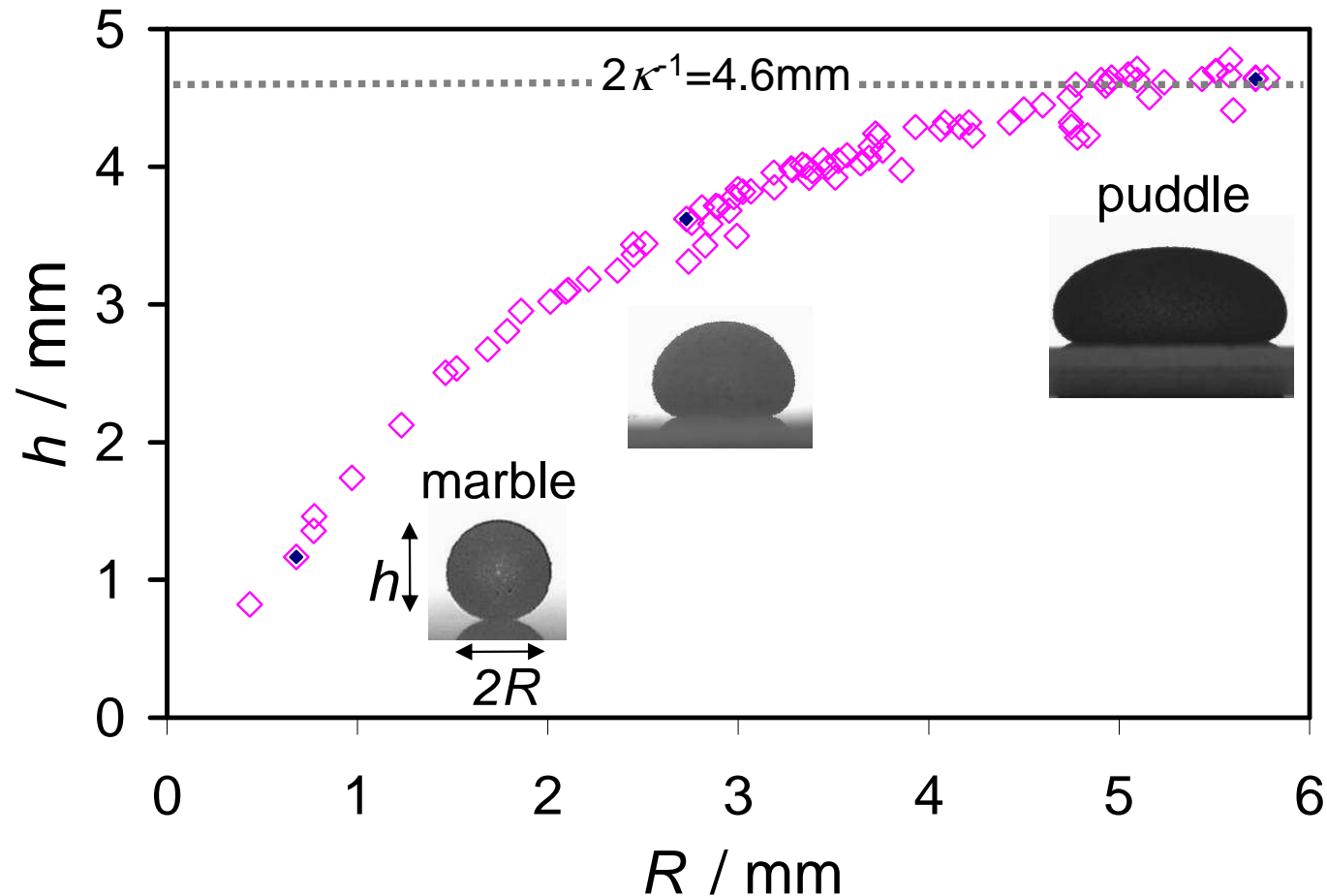
Hydrophobic lycopodium marbles are highly mobile  
Gravity distorts shape from spherical as volume increases

# Liquid Marbles – Size Data

Hydrophobic lycopodium plus 0.01M KCl deionised water

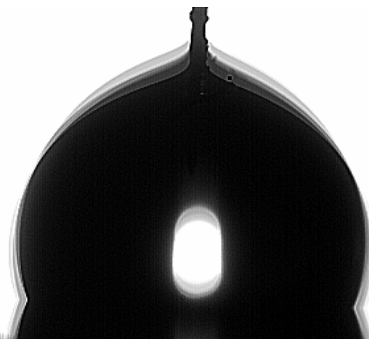
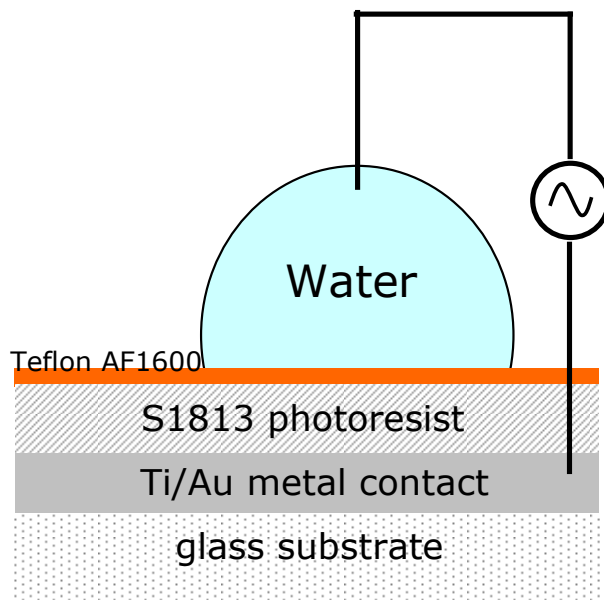
Transition from marble to puddle with increasing volume

Limiting value of puddle height gives twice the capillary length,  $h=2\kappa^{-1}$

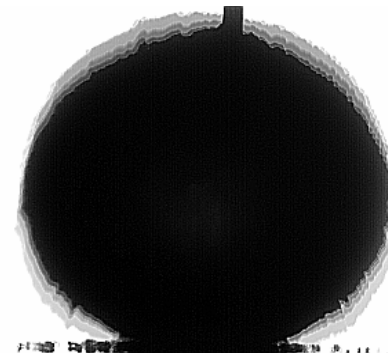
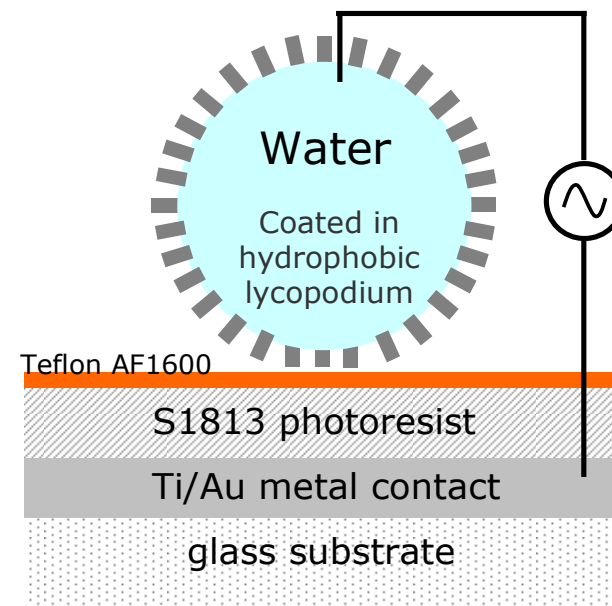


# Oscillation Experiments

# Experimental Arrangements



0.01M KCl,  $\sim 5 \mu\text{l}$ , 100 V, 44 Hz



0.01M KCl,  $5 \mu\text{l}$ , 100 V, 39 Hz

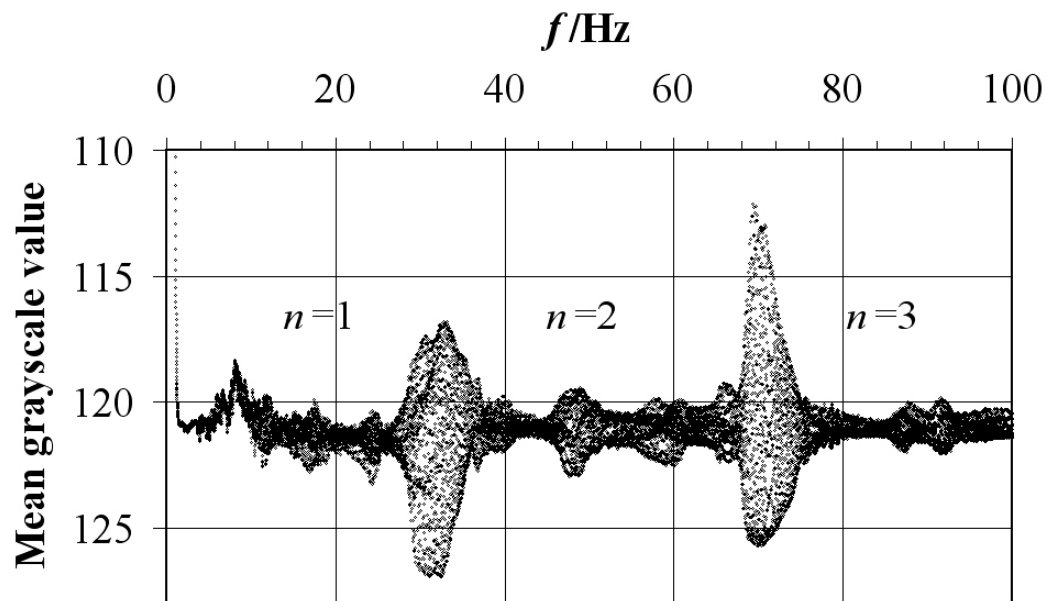
Hydrophobic lycopodium ( $17 \pm 3$ )  $\mu\text{m}$ , 5-275  $\mu\text{l}$  droplets of 0.01M KCl deionised water. Up to 200V p-p, 1 Hz-250 Hz sweep. High speed camera



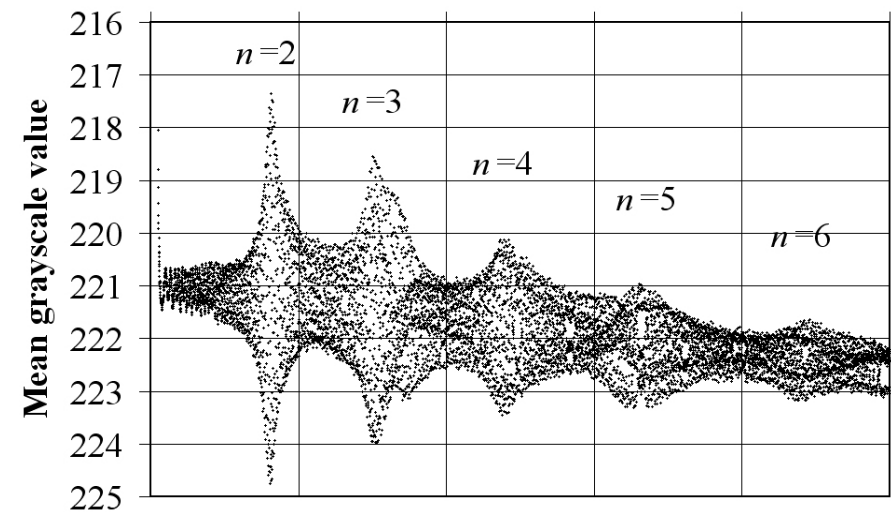
# Example Data

Example of the z-axis profile data used to locate resonant frequencies (0-100Hz)  
Mean grayscale value within a rectangular box selection at an anti-node close to the electrode wire for 100  $\mu\text{l}$  volumes

## Droplet

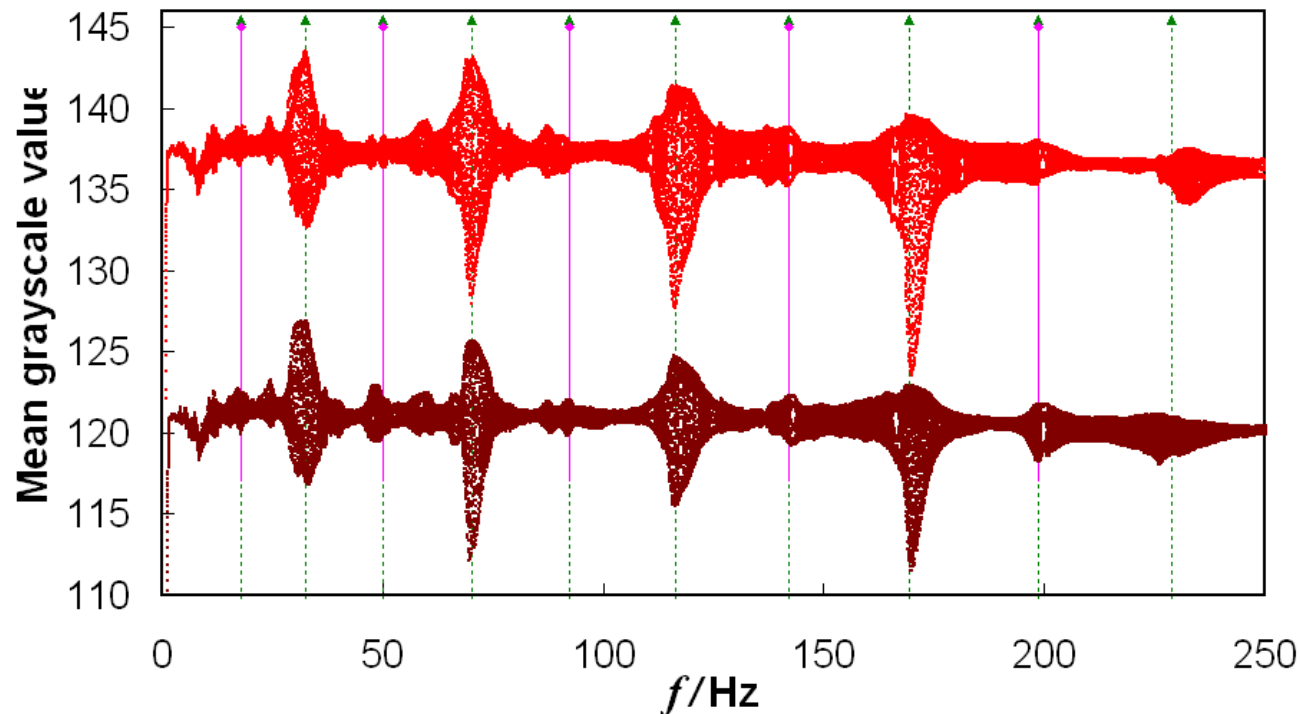


## Liquid Marble



# Analysis of 100 $\mu\text{l}$ Sessile Droplet

Fits using  $\rho=997 \text{ kg m}^{-3}$ ,  $\gamma_{LV}=72.8 \text{ mN m}^{-1}$ . Single fitting parameter used is  $\theta_e$  to reduce perimeter length,  $p$ , by 7% (modified from  $115^\circ$  to  $87^\circ$ ) to achieve a fit



Data from  
RHS of wire

Data from  
LHS of wire

Immobile contact line case: green ( $j=1, 2, \dots, 10$ ) (capillary only with  $\tanh()=1$ )

Mobile contact line case: pink ( $k=1, 2, \dots, 5$ ) (capillary only with  $\tanh()=1$ )

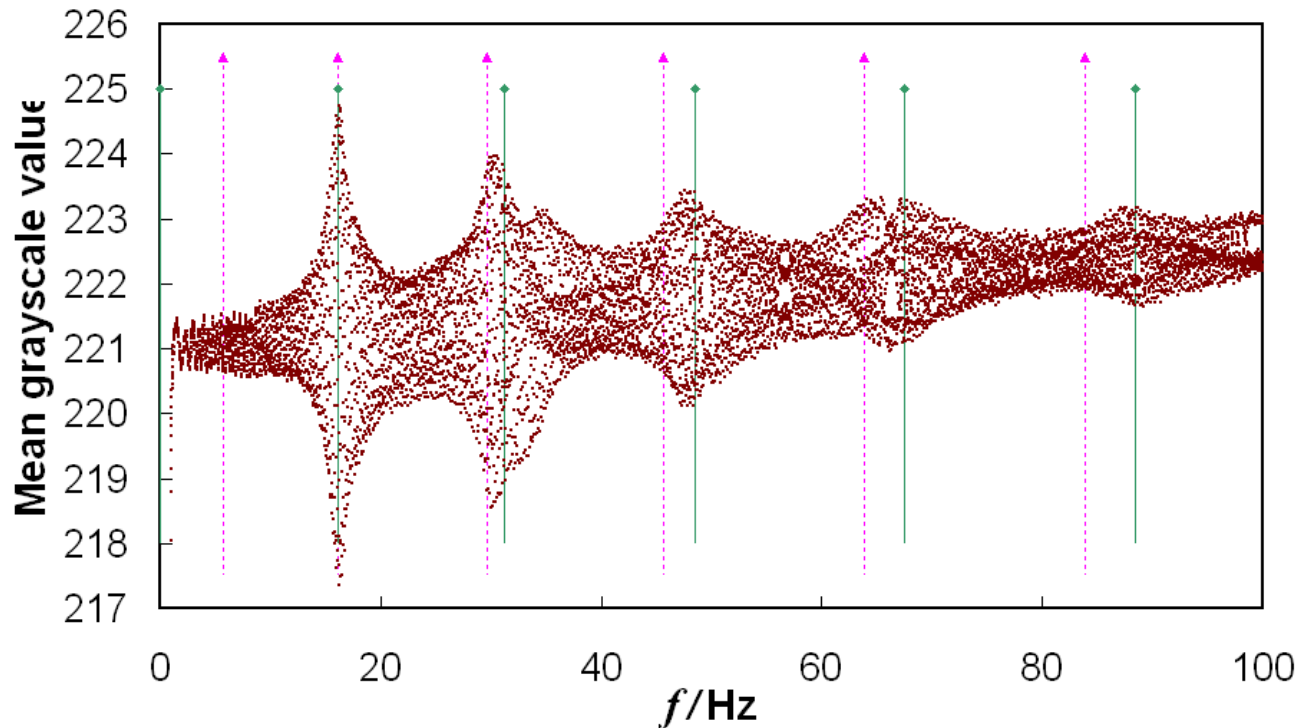
Major resonances occur for immobile contact line case with mode number  $j$  even

Resonances with  $j$  odd are reduced. These correspond to contact line having either node or anti-node, e.g. immobile  $j=3$  and mobile  $k=2$  both give  $2\lambda$  perimeter length and same  $f$

# Preliminary Analysis of 100 $\mu\text{l}$ Liquid Marble

Fits using,  $\gamma_{LV}=52 \text{ mN m}^{-1}$  determined from spherical to puddle data

Single fitting parameter used  $\rho=1690 \text{ kg m}^{-3}$  assuming spherical shape



Data from  
LHS of wire

Mobile contact line case: pink ( $j=1, 2, \dots, 6$ ) (capillary only with  $\tanh()=1$ )

Spherical droplet case: green ( $n=2, \dots, 6$ )

First major resonances occurs for  $n=2$  and  $j=2$  mode numbers

Spherical case using  $n(n-1)(n+2)$  formula appears excellent provided an effective density is used - ratio of surface tension to effective density gives a "spring constant"

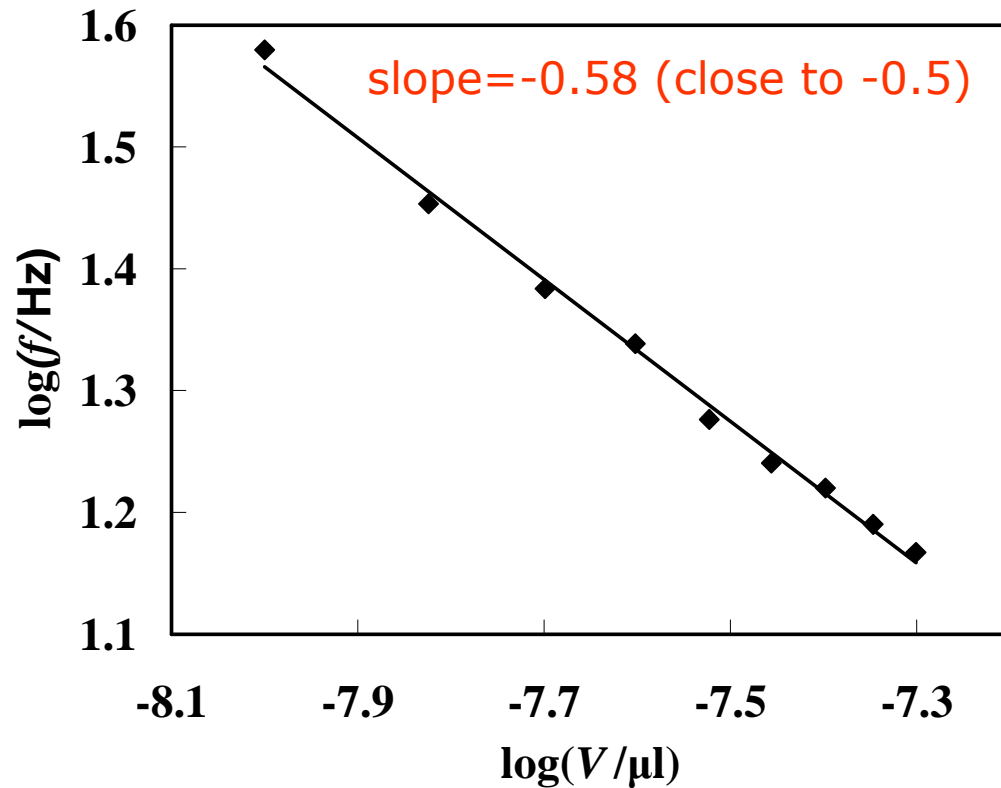
# Volume Dependence (Marble)

Assuming capillary dominated

$$f_k^2 \approx \frac{2\pi k^3 \gamma_{LV}}{p^3 \rho} \tanh\left(\frac{2\pi h k}{p}\right) \propto \frac{k^3}{p^3} \propto \frac{k^3}{V}$$

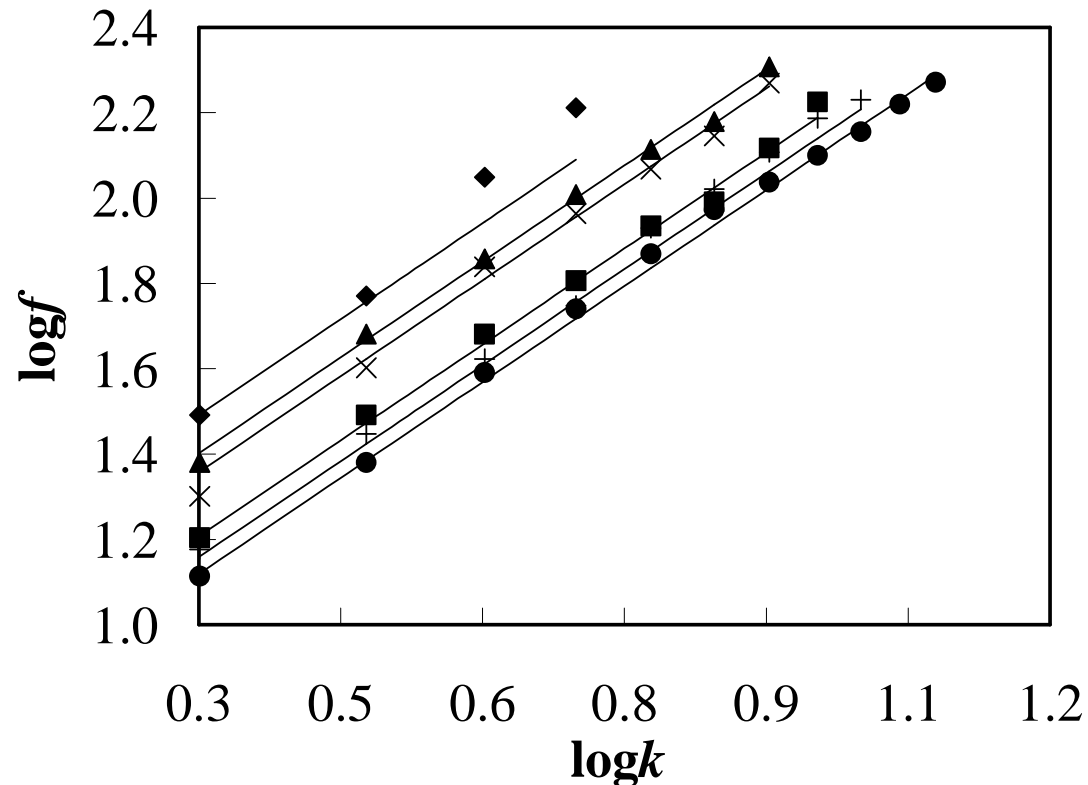
*Square of frequency as a function of reciprocal of marble volume*

Data for  $k=2$  mode ( $V=10-50 \mu\text{l}$ )



# Mode Dependence (Marble)

Log-log plot of frequency as a function of mode number for liquid marbles of volumes 10  $\mu\text{l}$  ( $\blacklozenge$ ), 30  $\mu\text{l}$  ( $\blacktriangle$ ), 50  $\mu\text{l}$  ( $\times$ ), 100  $\mu\text{l}$  ( $\blacksquare$ ), 125  $\mu\text{l}$  ( $+$ ), 150  $\mu\text{l}$  ( $\bullet$ ).



Solid lines are fits using 52  $\text{mN m}^{-1}$  ( $\rho=1690 \text{ kg m}^{-3}$  for 50, 100, 125 and 150  $\mu\text{l}$ ;  $\rho=2300 \text{ kg m}^{-3}$  for 30  $\mu\text{l}$ ;  $\rho=4560 \text{ kg m}^{-3}$  for 10  $\mu\text{l}$  (mode misidentification likely for 10  $\mu\text{l}$ ))

Fitting model: Capillary only with mobile contact line and  $\tanh(\cdot)=1$

Note: Spherical model can be fitted to data using lightly adjusted density values (spring constant)

# Summary

- Shape oscillations can be induced using an EWOD approach
- Droplet's main resonances for immobile contact line and  $j$  even
- Liquid marbles provide an idealized perfectly non-wetting system
- No need for active levitation/microgravity
- Zero contact angle hysteresis/mobile contact line
- Potential to study capillary-to-gravity regime transition

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The End

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See Also Our Poster  
*Microfluidic Switchable Diffraction Grating*  
*Gary G. Wells, **Carl V. Brown**, Glen McHale,  
Michael I. Newton and Christophe L. Trabi*

# See Poster - Microfluidic Switchable Diffraction grating

Gary G. Wells, Carl V. Brown, Glen McHale, Michael I. Newton and  
Christophe L. Trabi

School of Science & Technology, Nottingham Trent University, UK

6<sup>th</sup> International Meeting on Electrowetting

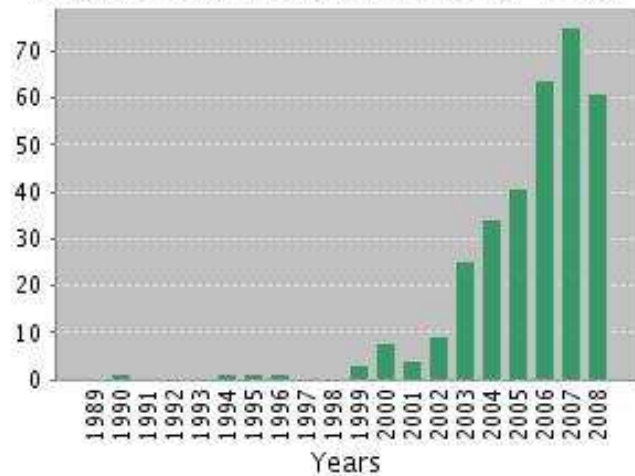
UCLA 20<sup>th</sup>-22<sup>nd</sup> August 2008

# Appendices

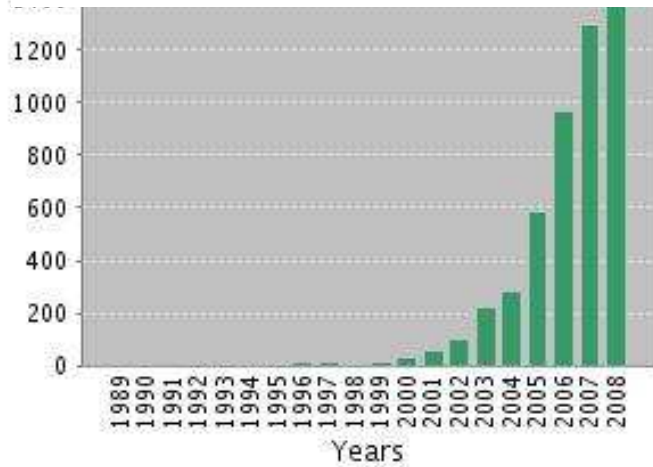
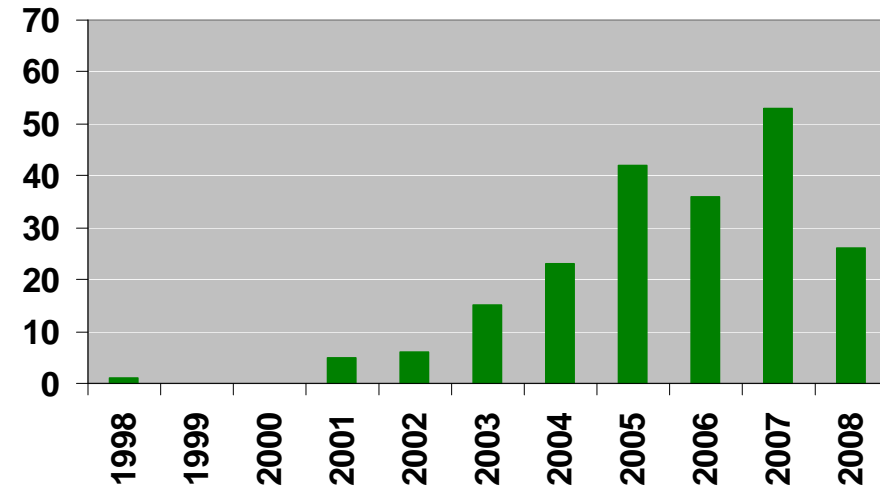


# “Electrowetting”

Published Items in Each Year



ISI Proceedings Papers in Each Year



# Fitting Equations

## Type 1 – Immobile Contact Line Analysis ( $k=2,3,4, \dots$ )

$$f_j^2 = \frac{(j+1)g}{4\pi p} \left( 1 + \left( \frac{(j+1)\pi}{p\kappa} \right)^2 \right) \tanh \left( \frac{(j+1)\pi h}{p} \right) \approx \frac{\pi(j+1)^3 \gamma_{LV}}{4p^3 \rho}$$

Sessile Droplet (side-view) Perimeter: side – view perimeter =  $2R_{cap} \theta$

$$\beta(\theta) = 2 - 3 \cos \theta + \cos^3 \theta \qquad R_{cap} = \left( \frac{3V}{\pi\beta} \right)^{1/3}$$

## Type 2 – Mobile Contact Line Analysis ( $k=2,3,4, \dots$ )

$$f_k^2 = \frac{kg}{2\pi\rho} \left( 1 + \left( \frac{2\pi k}{p\kappa} \right)^2 \right) \tanh\left( \frac{2\pi h k}{p} \right) \approx \frac{2\pi k^3 \gamma_{LV}}{p^3 \rho}$$

## Type 2 – Spherical Droplet Analysis ( $k=2,3,4, \dots$ )

$$f_n^2 = \frac{n(n-1)(n+2)\gamma_{LV}}{4\rho\pi^2 R^3} = \frac{2\pi n(n-1)(n+2)\gamma_{LV}}{p^3 \rho}$$

Spherical Droplet (side-view) Perimeter:  $p = 2\pi R = 2\pi \left( \frac{3V}{4\pi} \right)^{1/3}$

Sessile Droplet (side-view) Perimeter: side – view perimeter =  $2R_{cap} \theta$

$$\beta(\theta) = 2 - 3\cos\theta + \cos^3\theta \quad R_{cap} = \left( \frac{3V}{\pi\beta} \right)^{1/3}$$

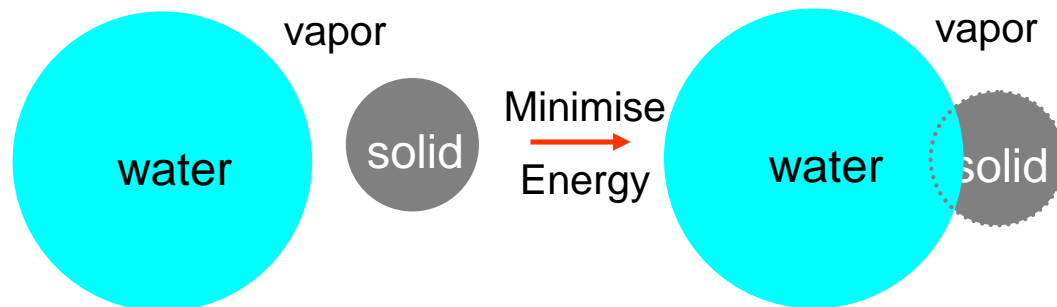
# Spare Slides

# Liquid Marbles

## Loose Surfaces

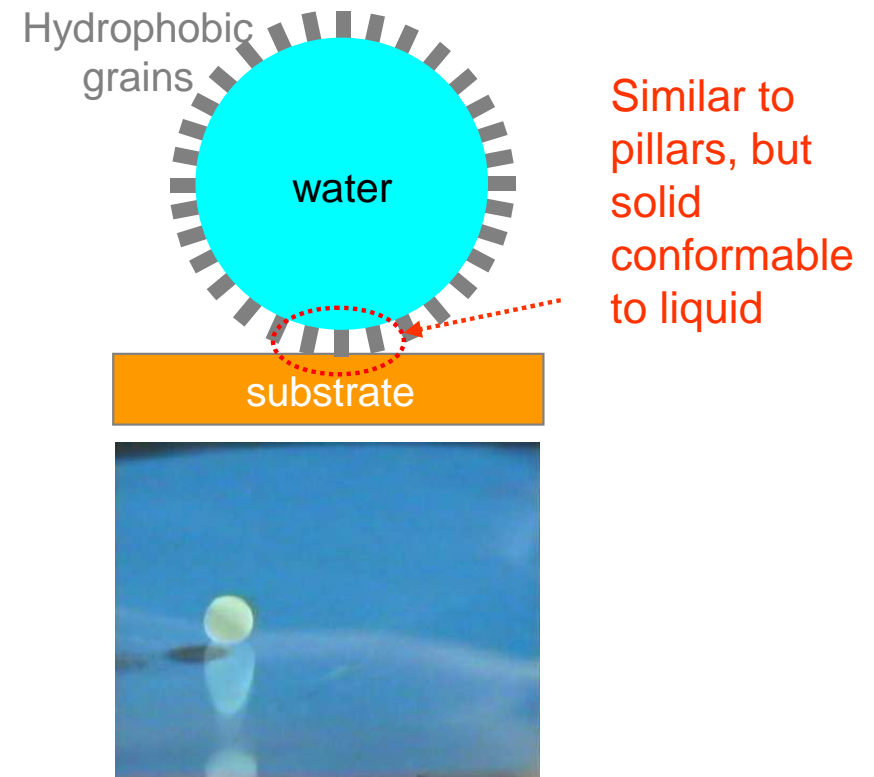
1. Grains are not fixed, but can be lifted
2. Surface free energy favors solid grains attaching to liquid-vapor interface
3. A water droplet rolling on a hydrophobic lycopodium (or other grain/powder) becomes coated and forms a liquid marble

## Hydrophobic Grains and Water



$$\Delta F = -\pi R_g^2 \gamma_{LV} (1 + \cos \theta_e)^2$$

Energy is always reduced on grain attachment

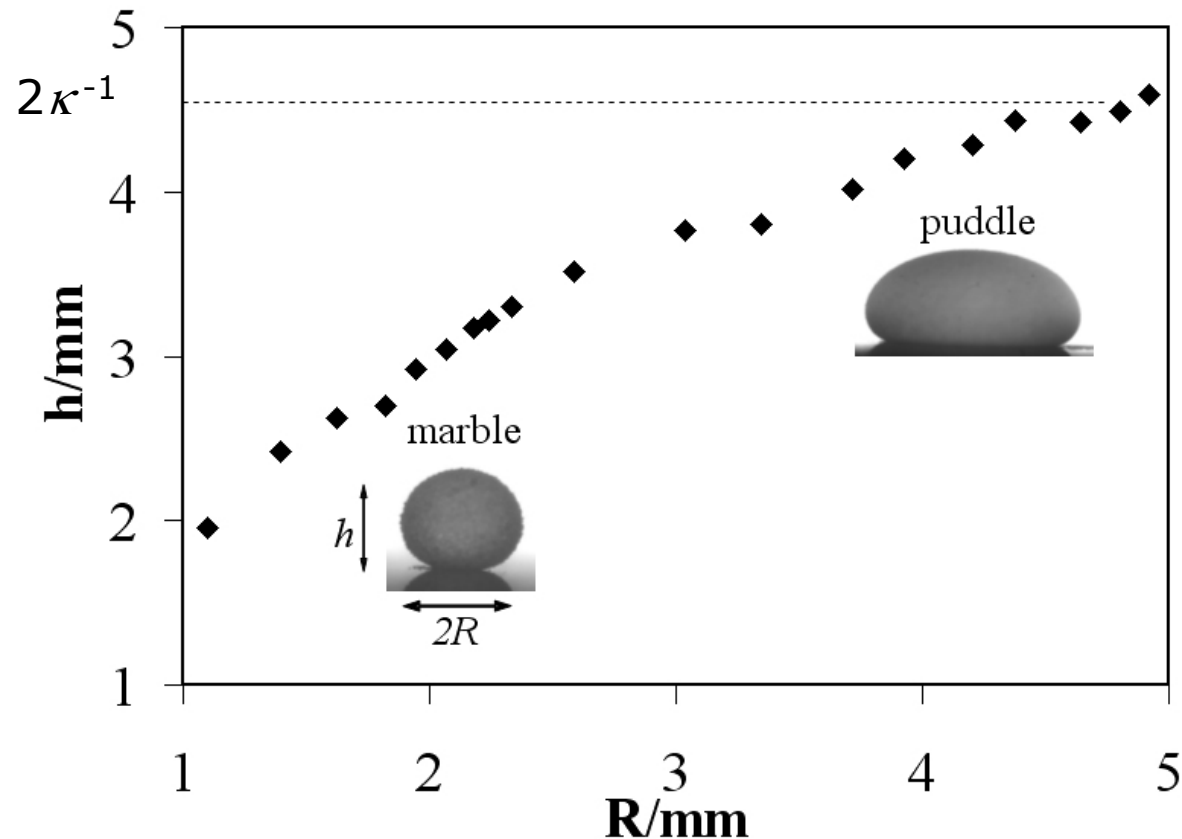


# Liquid Marbles – Size Data

Hydrophobic lycopodium plus 0.01M KCl deionised water

Transition from marble to puddle with increasing volume

Limiting value of puddle height gives twice the capillary length,  $h=2\kappa^{-1}$



Inset are images taken for marble of radius 1.1 mm and puddle of radius 4.8 mm.

# Mode Dependence (Droplets and Marbles – Subject to possible errors in droplet data)

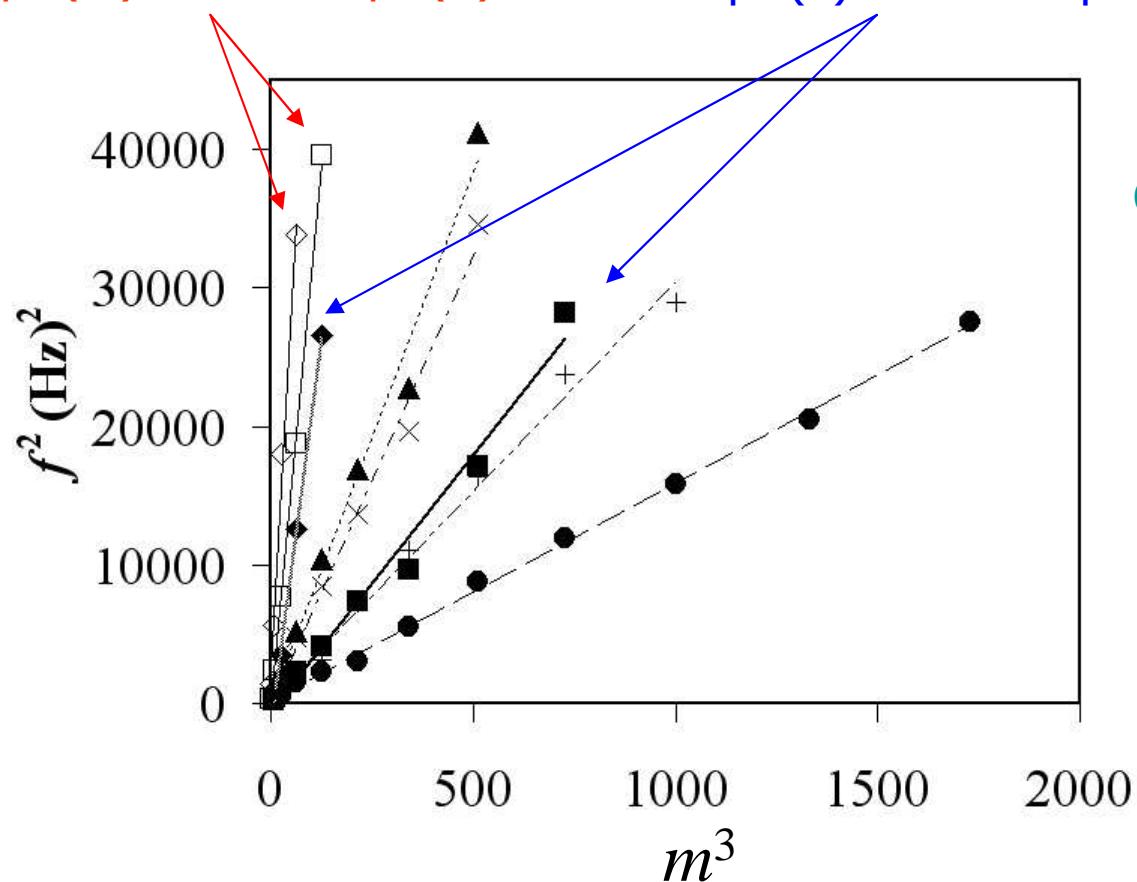
Square of frequency as a function of mode number,  $m$ , cubed

Droplets of volumes

10  $\mu\text{l}$  ( $\diamond$ ) and 100  $\mu\text{l}$  ( $\square$ )

Liquid marbles of volumes

10  $\mu\text{l}$  ( $\blacklozenge$ ) and 100  $\mu\text{l}$  ( $\blacksquare$ )



Note: Lines are trend lines and not fits